Numerical calculation methods for neutron noise in heterogeneous reactor cores

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- **6** Calculation of neutron noise in FR

- Neutron noise = the difference between the time-dependent flux and its mean value, provides interesting dynamic information of the core status and can be measured by in-core and ex-core detectors.
- The noise technique can be used online without disruption of the reactor operation.
- Possibility to investigate the neutron noise behaviour and interpret the measurement.
- Support the instrumentation development.
- Numerical simulation remains a challenge for describing detector signals and improving core surveillance.

2. Noise equation in multi-group diffusion theory

- To solve the neutron noise equation, it is necessary to define a noise source through the fluctuations of static cross sections, and the static state solutions of the *k_{eff}* and fluxes. Thus, the solution of the static equation is required.
- The static equation in multi-group diffusion theory:

$$-\nabla \cdot [D_{g} \nabla \phi_{g} (\mathbf{r})] + \Sigma_{t,g} \phi_{g} (\mathbf{r}) = \frac{1}{k_{eff}} \chi_{g} \sum_{g'} \nu \Sigma_{f,g'} \phi_{g'} (\mathbf{r}) + \sum_{g' \neq g} \Sigma_{s,g' \to g} \phi_{g'} (\mathbf{r})$$
(1)

where,

$$\chi_{g} = (1 - \beta) \chi_{g}^{p} + \sum_{j} \beta_{j} \chi_{g,j}^{d}.$$
 (2)

All symbols have their usual meaning.

The multi-group noise equation is derived from the space- and time-dependent diffusion equations by assuming that all time-dependent terms, $X(\mathbf{r}, t)$, can be split into a stationary component plus a small fluctuation as

$$X(\mathbf{r},t) = X_0(\mathbf{r}) + \delta X(\mathbf{r},t), \qquad (3)$$

By assuming that the fluctuations are small so that only the first order noise needs to be taken into account. Performing a Fourier transform of time-dependent terms, $\delta X(\mathbf{r}, t)$, as

$$\delta X(\mathbf{r},\omega) = \int_{-\infty}^{\infty} \delta X(\mathbf{r},t) e^{-i\omega t} dt.$$
(4)

2. Noise equation in multigroup diffusion theory

The first order neutron noise in multi-group diffusion theory is:

$$-\nabla \cdot [D_{g} \nabla \delta \phi_{g}(\mathbf{r}, \omega)] + \Sigma_{t,g}(\omega) \,\delta \phi_{g}(\mathbf{r}, \omega) = \frac{1}{k_{eff}} \chi_{g}(\omega) \sum_{g'} \nu \Sigma_{f,g'} \delta \phi_{g'}(\mathbf{r}, \omega)$$

$$+\sum_{g'\neq g} \sum_{s,g'\rightarrow g} \delta \phi_{g'}\left(\mathbf{r},\omega\right) + S_g\left(\mathbf{r},\omega\right),\tag{5}$$

where, $\delta \phi_g(\mathbf{r}, \omega)$ is the noise in group g. The frequency-dependent total cross section is

$$\Sigma_{t,g}(\omega) = \Sigma_{a,g}(\omega) + \sum_{g' \neq g} \Sigma_{s,g \to g'},$$
(6)
$$\Sigma_{a,g}(\omega) = \Sigma_{a,g} + \frac{i\omega}{v_g}.$$
(7)

 $\chi_g(\omega)$ is the frequency-dependent fission energy spectrum:

$$\chi_{g}(\omega) = \chi_{g} - \sum_{j} \chi_{g,j}^{d} \frac{i\omega\beta_{j}}{\lambda_{j} + i\omega}.$$
(8)

The noise source is calculated from the fluctuations of macroscopic cross sections

$$S_{g}(\mathbf{r},\omega) = -\delta\Sigma_{a,g}(\omega)\phi_{g}(\mathbf{r}) - \sum_{g'\neq g}\delta\Sigma_{s,g\rightarrow g'}(\omega)\phi_{g}(\mathbf{r}) + \sum_{g'\neq g}\delta\Sigma_{s,g'\rightarrow g}(\omega)\phi_{g'}$$
$$+ \frac{1}{k_{eff}}\chi_{g}(\omega)\sum_{g'}\delta\left[\nu\Sigma_{f,g'}(\omega)\right]\phi_{g'}(\mathbf{r}).$$
(9)

3. Spatial discretization

The balance (static or noise) equation in each group in a node n is

$$\sum_{n'=1}^{N} A_{n,n'} J_{n,n'} + \Sigma_t \varphi_n V = Q_n V$$
(10)

 $J_{n,n'}$: the surface-averaged net current from node *n* through the interface with a neighboring node n',

- $A_{n,n'}$: the area of the interface of the two nodes,
- φ_n : the static flux or the noise,
- Q_n : the total source.



1 Nodal methods

- Advantages: can handle a large node size with correction of the net current.
- **Disadvantages:** not applicable in the cases of point-like sources and/or vibration.



1 Finite difference (FD) approximation

- **Advantages:** simple and flexible with node size, suitable to the cases of point-like sources and/or vibration.
- **Disadvantages:** problem size may be large, slow conversion speed for the cases of point-like sources and/or vibration



1 Finite difference (FD) approximation

The net current $J_{n,n'}$ is calculated as

$$J_{n,n'} = a_{n,n'} \left(\varphi_n - \varphi_{n'} \right). \tag{11}$$

The coupling coefficient $a_{n,n'}$ is

$$a_{n,n'} = \frac{2}{h} \frac{D_n D'_n}{D_n + D'_n},$$
(12)

If the neighboring node n' is a boundary, the current is

$$J_{n,n'} = \frac{(1-\alpha)2D_n}{h+4D_n}\varphi_n \tag{13}$$

where, a vacuum boundary has $\alpha = 0$ and a reflective boundary has $\alpha = 1$.

Depending on the size of problem $(N \times G)$ to chose appropriate solution methods.

1 Green/Transfer function

$$\delta\phi(\mathbf{r},\omega) = [G_{XS}(\mathbf{r},\mathbf{r}',\omega) \times \phi(\mathbf{r}']\delta_{XS}(\mathbf{r}',\omega)$$
(14)

- Advantages: Transfer function is solved only one, and applicable to any case of noise calculation with any source type.
- **Disadvantages:** time and memory consuming, not applicable to a large case with multi-energy groups

4. Numerical methods

- Finite difference (FD) approximation
- Comparison with analytic solution in case of point-like source



Figure 1: Comparison of fast and thermal noise in a PWR between numerical and analytic calculations.

- 2D, two-groups model
- Inite difference for spatial discretization
- Iner meshes for vibrating assembly
- Transfer function



Figure 2: Magnitude and phase of thermal noise induced by vibration of assembly M10 in x-direction (f = 8Hz).

Space-dependent noise induced by fuel assembly vibration



Figure 3: Magnitude and phase of thermal noise induced by trajectory vibration of assembly M10 in xy-direction (f = 8 Hz).

Ex-core noise induced by fuel assembly vibration

- The increase of noise amplitude with burnup is only found for a subset of fuel assemblies located at the periphery.



Figure 4: APSD (a.u.) of ex-core detector induced by simultaneous stochastic vibrations of fuel assemblies J9, K11 and L14 in a PWR.

Ex-core noise induced by fuel assembly vibration

- The noise induced by the peripheral assemblies dominates the ex-core detector signals.



Figure 5: APSD (a.u.) of detectors N1 and N2 induced by simultaneous stochastic vibrations of assemblies B5, E9, K6 and L14 in a PWR.

- Two modules: static and noise
- Ø Multigroups, several delayed neutron precursors
- Oiffusion theory, hexagonal geometry, finite differences, power iteration
- OMFD acceleration



Figure 6: Triangular discretization and flowchart of noise solution for hexagonal

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- SFR core
- 33 energy groups, 6 delayed neutron precursors
- One source: perturbation of coolant density at the core central



Figure 7: Triangular discretization and flowchart of noise solution for hexagonal reactor.

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 Energy-dependent noise induced by perturbation of coolant density at the core central.



Figure 8: Amplitude and phase of noise at different locations in the core: Source position, near source, far from source and Reflector (f = 1.0 Hz).

 Space-dependent noise induced by perturbation of coolant density at the core central.



Figure 9: Amplitude and phase of space-dependent noise in fast (group 9) and epi-thermal (group 20) groups (f = 1Hz).

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